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# Technical Note 71

# CALIBRATION OF FIVE GAMMA-EMITTING NUCLIDES FOR EMISSION RATE

U. S. DEPARTMENT OF COMMERCE
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*Technical Note*

71

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# CALIBRATION OF FIVE GAMMA-EMITTING NUCLIDES FOR EMISSION RATE

J. M. R. Hutchinson

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ABSTRACT

Mercury-203 and niobium-95 were calibrated by a $4\pi\beta$-$\gamma$ coincidence method for $\gamma$-emission rate, zinc-65 by comparison with the 1.12 Mev peak of scandium-46, sodium-22 by a $\gamma$-annihilation-quanta coincidence method and by a triple coincidence method, and strontium-85 by x-$\gamma$ coincidence counting. The accuracy of the calibration in all cases was $\pm 2$ %. The half-life of the isomeric state of rubidium-85 was measured and found to be .98 microseconds.

# Calibration of Five $\gamma$-Emitting Nuclides for Emission Rate

J. M. R. Hutchinson

## Introduction

In order to supply a need for point source $\gamma$-ray standards for calibration purposes in nuclear physics, five $\gamma$-emitting nuclides have been calibrated by the coincidence method for $\gamma$-emission rate. The nuclides are listed below with some of their properties.

Table I

| Nuclide | Emitted rays | $\gamma$ Energy(mev) | Half-life |
|---|---|---|---|
| Mercury-203 | $\gamma$, $\beta^-$ | 0.245 | 47d |
| Strontium-85 | $\gamma$, x (k capture) | 0.513 | 64d |
| Niobium-95 | $\gamma$, $\beta^-$ | 0.765 | 35d |
| Zinc-65 | $\gamma$, $\beta^+$, x | 1.119 | 245d |
| Sodium-22 | $\gamma$, $\beta^+$ | 1.28 | 2.8y |

The nuclides will be provided in solid, point source (diameter of deposit $<$ 1 cm) standard samples of $\gamma$-emission rate. Ideally one desires nuclides emitting a single $\gamma$-ray, whose energies are well known and evenly spaced with respect to energy from 0.1 Mev to 2.0 Mev and whose half lives are long enough to make them useful as accurate standards for a period of years. In the kit initially produced the range of

energies used is indicated in Table I. The activity of each standard is approximately $5 \times 10^4$ $\gamma$'s/sec. The estimated accuracy of the standardization for $\gamma$-emission rate is ± 2%.

This article will be divided into three parts devoted to (i) methods of source preparation, (ii) a description of the calibration of the standards, (iii) various checks on the $\gamma$-emission rate of the standards.

## SOURCE PREPARATION

### Mounting

It is desirable that in a small solid-angle, at least, the $\gamma$ radiation emitted by the source should be unscattered. Also the mounting has to be firm in order to allow a fixed geometry intercomparison with other sources of the same nuclide. Toughness and resistance to temperature changes are also desirable qualities. All these requirements were satisfied by covering an annulus of aluminum 2" o.d. 1 1/2" i.d. with thin adhesive polyester tape under tension. The source was deposited on the adhesive side of the tape, approximately in the center and covered with another layer of similar polyester tape.

Source material was received in highly concentrated form and diluted to a strength of approximately $5 \times 10^4$ $\gamma$'s/sec per 50 mg. of water solution. 5 ml. ampoules containing accurately weighed amounts of the diluted solution were prepared and sealed. The ampoules were

then measured in the N.B.S. ionization chamber for ion current per gram of solution relative to a fixed radium source (See Emission Rate Checks). The ampoules were then broken open and a number of sources of approximately 50 mg were weighed onto the source mounts using a picnometer and a balance accurate to 0.1 mg. The latter operation was performed quickly to prevent much evaporation during the weighing. The accuracy of these weighings in the worst case as checked by relative counting rates in a fixed geometry was found to be ± 0.6%. Since each source was checked for counting rate relative to a fixed standard, the error in intercomparing standard samples is less than ± 0.4%. The sources were allowed to dry, after which tape was placed over them securing them in a sandwich.

## CALIBRATION

### Mercury-203 ($4\pi\beta$-$\gamma$ coincidence counting)

The energy of the conversion electron near the high energy limit of the continuous $\beta$ spectrum caused difficulty in monitoring $\beta$'s in the continuous spectrum without counting the conversion electrons. Moreover the total conversion is too large (22.6%) to be corrected for accurately in any easy manner in low efficiency counters. The obvious solution is to use a $4\pi\beta$ counter which has high enough efficiency so that most of the time a beta is counted in, say a coincidence apparatus, a conversion electron is also counted, the two being recorded as one pulse in the $\beta$ channel. For 100%

β efficiency the counts in the β channel equal the β emission rate alone without any contribution from the accompanying conversion electrons.

In view of these considerations a $4\pi\beta$-$\gamma$ coincidence measurement of $Hg^{203}$ was made. Since the coincidence analyser could not accept counting rates of $10^4$/sec an accurate further dilution of the source material was made and deposited by picnometer onto thin collodion source holders. Several sources were counted for β's, γ's and coincidences. Approximate disintegration rates and efficiencies were determined after correcting for background in all channels. The observed β efficiency was approximately 94%. A final disintegration rate was determined after correcting for (i) accidentals (ii) dead time losses (iii) internal[1] conversion. The accidental rate was given by $N_\beta N_\gamma \tau(2-\epsilon_\beta-\epsilon_\gamma)$. The dead time loss rate for both channels was small.

A multiplicative factor correcting for internal conversion is

$$1 + \frac{\alpha(1-\epsilon_\beta)\epsilon_{ce}}{\epsilon_\beta(1+\alpha)}$$

Thus the final disintegration rate is given by:

$$N_o \quad \frac{(N_\beta - B_\beta)(N_\gamma - B_\gamma)}{(N_c - B_c - A)} \quad \frac{1}{1 + \frac{\alpha(1-\epsilon_\beta)\epsilon_{ce}}{\epsilon_\beta(1+\alpha)}}$$

$N_\beta$ = β channel rate

$N_\gamma$ = γ channel rate

$\epsilon_\beta$ = β efficiency

$\epsilon_\gamma$ = γ efficiency

$\alpha$ = total conversion coefficient

$\epsilon_{ce}$ = conversion electron efficiency

$\tau$ = resolving time ( $1.45 \times 10^{-6}$ sec)

and inserting some typical numbers (all in per second)

$$N_o = \frac{(356.87-0.90)(11.879-4.109)}{(7.3227-0.1160-0.00639)} \quad \frac{1}{1.0113}$$

$\epsilon_{ce}$ was taken as unity

$N_c$ = coincidence rate

$N_o$ = disintegration rate

$B_\beta$ = $\beta$ background

$B_\gamma$ = $\gamma$ background

$B_c$ = coincidence background

A = accidental rate

A method similar to the above was used to calibrate $Nb^{95}$ and to check the disintegration rate of $Na^{22}$ described in later sections of this article. A total conversion coefficient of 22.57% based on the numbers of Nijgh and Wapstra[2] was used.

Strontium-85 (x-$\gamma$ coincidence count)

The long half-life ($10^{-6}$ sec) of the isomeric state of $Rb^{85}$ required special consideration when making an x-$\gamma$ coincidence count on the decay of $Sr^{85}$. In order that all the coincidences be recorded, a resolving time long compared to the half life of the isomeric state was necessary. It was also important to make the resolving time small enough to minimize the error due to variation in accidental rate. The half-life was determined and coincidence counts were made at two resolving times.

Half-life determination of the 0.513 Mev state of $Rb^{85}$

The schematic experimental arrangement used in the half-life determination is shown in Figure IV. A 1-mm thick NaI crystal with

an electron-multiplier phototube was used to detect the x-rays and a 2" x 2" NaI crystal with phototube was used to detect the 0.513-Mev $\gamma$ radiation. The $\gamma$ ray was received by the $\gamma$ channel some time after the x-ray had been detected in the x-ray channel. To measure the half-life of the isomeric state a plot on semi-log paper was made of the number of real coincidences divided by the product of the two singles rates versus delay in the x channel for a resolving time short compared to the half life. A straight line resulted and the half-life was read from the graph. The resolving time was 0.2 μs. The accidentals were found by increasing the delay to 20 μs. The variable delay was calibrated with a crystal oscillator. Six values of the half life of the isomeric state of $Rb^{85}$ were obtained - 0.96, 0.96, 0.97, 1.00, 0.98, 1.00 giving an average corrected value of

$$\tau_{1/2} = 0.98 \pm 3\%$$

The major source of error in this measurement is due to inaccurate calibration of the time delay.

Calibration of disintegration rate of $Sr^{85}$

The disintegration rate of a $Sr^{85}$ standard sample was measured by coincidence counting with the equipment of Figure IV using resolving times of 2.57 μs and 8.5 μs in the coincidence analyser. Knowing the half life of the isomeric state of $Rb^{85}$ and the resolving time of the coincidence circuit, the fraction of the total real coincidences which were lost due to the delay of the $\gamma$ in decaying from the isomeric state

were calculated and a correction in the disintegration rate was made for them. To obtain the 8.5 μs resolving time a gate 7.3 μs long was placed in the γ channel and a pulse of 1.2 μs shaped by a shorted delay line was introduced into the x channel. The time relationship between the two channels was found by placing an annihilating positron source between the detectors and adjusting the delay in the x channel to where coincidences just disappeared, i.e., where the total delay in the x channel was 7.3 μs longer than the delay in the γ channel. The channels were then adjusted to overlap by approximately 0.6 μs (i.e., the delay in the x channel was reduced by 0.6 μs) to remove any loss of coincidences due to differing rise times of the pulses in various parts of the circuit. Under these circumstances γ rays which were emitted within 7.3 + 1.2 - 0.6 = 7.9 μs after the associated x-rays were emitted could produce a coincidence. The coincidence equation for this experimental arrangement is then:

$$N_o = \left(1 - \frac{1}{2\frac{7.9}{0.98}}\right) \frac{N_x N_\gamma}{N_c}$$

where $N_x$ = singles counts in the x channel, $N_\gamma$ = singles counts in the γ channel, $N_c$ = coincidences.

A coincidence count of $Sr^{85}$ was also made with resolving times shaped by a shorted delay line of $\frac{2.57}{2}$ μs in both channels. The pulses were made to overlap by 1.034 μs. The disintegration rate for this setup was then calculated with the following formula

$$N_o = \left(1 - \frac{1}{2\frac{4.11}{0.98}}\right) \frac{N_x N_\gamma}{N_c}$$ where 4.11 = 2.57 - 1.034 + 2.57.

The disintegration rate of the $Sr^{85}$ master solution obtained using these two resolving times agreed with each other within 1% and both fell within 1% of the predicted $N_o$/gm of the master solution using the ionization chamber curve.

Zinc-65 (comparison with the 1.12 mev scandium-46 peak)

Zinc was calibrated for $\gamma$ emission rate by comparing the 1.112 Mev[3] peak of $Zn^{65}$ as seen in a 2" x 2" crystal with the upper 1.119 Mev[4] peak of a calibrated $Sc^{46}$ source. The disintegration scheme of $Sc^{46}$ is similar to that of $Co^{60}$.

A $Sc^{46}$ source was calibrated by $\gamma$-$\gamma$ and $\beta$-$\gamma$ coincidence counting. The $\gamma$-$\gamma$ coincidence method used in calibrating $Sc^{46}$ followed that of Hayward, Hoppes and Mann in their calibration of $Co^{60}$.[5] Briefly, their method was to set their crystals at different angles with respect to the source and monitor both channel's singles rates and the coincidence rate. Corrections for the angular correlations of the cascade $\gamma$'s and the finite size of the crystals can be computed using theoretical values given by Rose.[6] The fact that the crystals cannot be made to monitor the $\gamma$'s separately can be taken into account by adjusting the discriminator settings on the amplifier outputs to give equal efficiencies and by introducing a factor of 1/2 into the coincidence equations.[7] For two, two inch crystals each at a distance of four inches from the source, the theoretical ratio $\frac{N_1 N_2}{N_c}$ at $\pi$ and $\frac{\pi}{2}$ for a 4-2-0 transition is 1.155. The experimental ratio was 1.155 ± 0.010.

A $\beta$-$\gamma$ measurement on the same source was also made and the value for the disintegration rate fell within 0.5% of that obtained by $\gamma$-$\gamma$ coincidence counting. The estimated accuracy of the measurement of the $Sc^{46}$ source was ± 0.5%.

A comparison of the $Sc^{46}$ and $Zn^{65}$ sources was made by comparing the counting rate in the 1.12 $Sc^{46}$ peak as "seen" through a 5v window with the counting rate of the $Zn^{65}$ $\gamma$ peak as seen through the same window. Measurements were made at different settings to check the effect on the ratio of $\gamma$ emission rates of the slightly differing energies of the two $\gamma$'s. Comparison in this manner of a $Sc^{46}$ source and a $Zn^{65}$ source provided the same ratio of $Sc^{46}$ to $Zn^{65}$ for differential discriminator settings varying over a range of 0.03 Mev - i.e., a range known to be greater than the difference in energy between the upper $Sc^{46}$ peak and the $Zn^{65}$ peak.

Sodium-22 (triple coincidence and $\gamma$-annihilation quanta coincidence)

Sodium-22 decays 90% by positron emission and 10% by electron capture to the 1.28 Mev excited state of $Ne^{22}$. This state then decays promptly to the ground state. The coincidence experiment was as follows.

A source was encased in a lead sheath to produce complete annihilation of the positrons within a small volume. The annihilation quanta were received in coincidence in $a_1$ and $a_2$ which were 180° apart with respect to the source (see Figure VIII) and fed through amplifiers into coincidence analyser $C_a$. The coincidence pulses out of $C_a$ represented a sampling of positrons. These pulses were fed into coincidence

analyser $C_{\beta\gamma}$. Pulses into the $\gamma$ channel were amplified ($A_\gamma$) and discriminated against annihilation quanta (D) and also fed into $C_{\beta\gamma}$. Thus at $C_{\beta\gamma}$ one had the arrangement for an oridnary $\beta$-$\gamma$ coincidence experiment. $\beta$, $\gamma$, and coincidence rates were taken. Accidental and background rates were also taken. Turning to the "$\beta$" channel again, part of the "$\beta$" counts resulted from coincidences between the 1.28 Mev $\gamma$ and annihilation photons. This so-called $\gamma$ background in the "$\beta$" channel and other effects such as accidentals in the "$\beta$" channel etc. were corrected for by subtracting the counts in $S_\beta$ when $a_1$ was rotated away from the 180° position with respect to the source from the "$\beta$" singles rate. The advantage of this method lies in the fact that the "$\gamma$ background in the $\beta$ channel" can be subtracted directly without making an independent estimate of the 1.28 Mev $\gamma$ efficiency in the 0.511 Mev region.

$\gamma$ - Annihilation $\gamma$ Coincidence Count

A coincidence count was then taken between the 1.28 Mev $\gamma$ and the annihilation $\gamma$'s. The experimental arrangement consisted of an ordinary two channel coincidence analyser with two, 2" x 2" crystals facing each other but which had the freedom to rotate with respect to the line connecting the source and other crystal. One channel was fed through a window which contained the annihilation peak. The positrons were again completely annihilated close to the source by a lead sandwich. The coincidence equations are in general:

$$N_\gamma = N_o \epsilon_\gamma + n^\gamma_{\beta\gamma}$$

$$N_\beta = N_o \epsilon_\beta + n'_\gamma + n^\beta_{\beta\gamma}$$

$$N_c = N_o \epsilon_\beta \epsilon_\gamma + n^\gamma_{\beta\gamma} \epsilon'_\beta$$

Nomenclature.

$n^\gamma_{\beta\gamma}$ = sum coincidence in $\gamma$ channel

$n^\beta_{\beta\gamma}$ = sum coincidence in $\beta$ channel

$n'_\gamma$ = 1.28 Mev comptons in $\beta$ channel

$\epsilon'_\beta$ = probability that a $\beta$ will be detected through $\beta$ window if the $\beta$ strikes the crystal

Hence if the two crystals oppose each other with the source in the center we get:

$$\frac{N_\gamma N_\beta}{N_c} = \frac{(N_o \epsilon_\gamma + n'_\gamma + n^\gamma_{\beta\gamma})(N_o \epsilon_\beta + n^\beta_{\beta\gamma})}{N_o \epsilon_\beta \epsilon_\gamma + n^\gamma_{\beta\gamma} \epsilon'_\beta}$$

If the crystals are not aligned with the source $n^\gamma_{\beta\gamma} \epsilon'_\beta = 0$ since the annihilation quanta are emitted in opposite directions. From a detailed analysis of the $Na^{22}$ spectrum $\epsilon'_\beta$ was estimated at 50%. $n'_\gamma$ was measured by counting the Compton recoils from a calibrated $Co^{60}$ source through the $\beta$ window and by assuming that the two $Co^{60}$ $\gamma$'s with an average energy of 1.25 Mev would produce twice as many Compton recoils per disintegration in the region of 0.511 Mev as one $Na^{22}$ $\gamma$ with an energy of 1.28 Mev.

$n^\gamma_{\beta\gamma} \epsilon'_\beta$ ( $0.5\, n^\gamma_{\beta\gamma}$) was estimated from the difference between $\frac{N_\beta N_\gamma}{N_c}$ when the crystals and the source were and were not aligned. Some typical numbers will help illustrate the calculation. $\frac{N_\beta N_\gamma}{N_c}$

at 180° was found to be 20% smaller than at 90°. Thus $N_{\beta\gamma}' \epsilon_\beta$ was approximately 20% of the true coincidence rate. $N_c$ was about 1 count/sec, hence $n_{\beta\gamma}^{\gamma} \epsilon_\beta' (0.5\ n_{\beta\gamma}')$ = (20%) 1 count/sec = 0.2 counts/sec. Hence $n_{\beta\gamma}^{\gamma}$ = 0.4 counts/sec. Now $N_\gamma$ was about 70 counts/sec hence the summing correction in the $\gamma$ singles channel was approximately $\frac{0.2}{(70)(0.5)}$ = 0.6%. With regard to the "β" channel, it was evident that the summing should be less than in the $\gamma$ channel because the "β" channel was "seen" through a window and the $\gamma$ channel was not. Analysis of the spectrum showed this estimate to be correct--in fact the summing in the β channel was calculated to be approximately 0.1% of the annihilation $\gamma$ counting rate.

$N_o$ was counted at an angle of 180°. After the above corrections were made and the background and accidental coincidences were subtracted, a disintegration rate of the $Na^{22}$ agreeing within 1% of the triple coincidence measurement was counted.

Two Other Checks on the Calibration of $Na^{22}$

Calibrated sources of $Sr^{85}$ and $Co^{60}$ were used in the right proportions to reproduce the spectrum of $Na^{22}$ in the region of the annihilation peak, the $Co^{60}$ reproducing the Compton-recoil background. Counting through a 6v window (this just took in the annihilation peak) an intercheck on the calibration of three sources was made since the $Co^{60}$ gave degraded quanta which were equivalent to the 1.28 Mev $Na^{22}$ quanta in that energy region and the $Sr^{85}$ produced $\gamma$'s which were of the same energy (within 0.5%) as the annihilation quanta. The following

formula was used to intercompare the sources:

$$\frac{N_s}{N_{os}} = 2(0.898)\frac{N_r}{N_{or}} + \frac{N_c}{N_{oc}},$$

where $N_s$, $N_r$, $N_c$ are the counting rates of $Na^{22}$, $Sr^{85}$ and $Co^{60}$ sources respectively and $N_{os}$, $N_{or}$, $N_{oc}$ are the corresponding disintegration rates.

In this equation the fact that two annihilation quanta appear for every positron in a $Na^{22}$ disintegration and the fact that for each decay of $Co^{60}$ there are two $\gamma$'s emitted were taken into account. 0.898[8] is the positron branching fraction from a decaying $Na^{22}$ nucleus. Using an average value of disintegration rate of both double and triple coincidence counting for $Na^{22}$, the left hand side of the equation agreed with the right hand side to within 1.5%.

A second intercomparison was provided by the fact that the peak efficiencies of the upper peak of $Co^{60}$, the 1.28 Mev peak of $Na^{22}$ and the peak of $Zn^{65}$ fell on a straight line on a log-log-graph in agreement with the curves of Lazar, Davis and Bell.[9]

EMISSION RATE CHECKS

In addition to the source calibration two checks on the accuracy of calibration of the disintegration rate of the standard samples were made, these were namely (i) the final weighed sources were intercompare for counting rate thus checking the weighed amount of the master solution, and (ii) a comparison with a group of nuclides calibrated by

$4\pi\beta$-$\gamma$ coincidence counting which included among others $Na^{22}$, $Nb^{95}$, $Hg^{203}$ but not $Sr^{85}$, $Zn^{65}$ was made using the ion chamber as a means of comparison. A smooth curve was drawn through the points given by these $4\pi\beta$-$\gamma$ calibrated nuclides on a graph of $\frac{\text{ion current}}{\text{gm } \gamma/\text{sec}}$ versus energy. Interpolated values of $\frac{\text{ion current}}{\text{gm } \gamma/\text{sec}}$ were obtained for $Sr^{85}$ and $Zn^{65}$, and a comparison of these nuclides was made. It should be emphasized that in the cases of $Sr^{85}$ and $Zn^{65}$ the quoted values for the $\gamma$ emission rates of the samples were those obtained directly by other methods, the ion chamber measurements served only to check the limits to within 2 or 3%. As it turned out, calibration of these two nuclides agreed to within ± 1% of the calibration of the same two nuclides using interpolated values on the $\frac{\text{ion current}}{\text{gm } \gamma/\text{sec}}$ versus energy curve. The disintegration rate of the $Na^{22}$ master solution obtained in other ways agreed with that obtained by $4\pi\beta$-$\gamma$ coincidence measurement.

In conclusion the author gratefully acknowledges continued interest in the work by R.W. Hayward, numerous helpful discussions with S. Garfinkel and D. D. Hoppes, and aid with the dilutions by R. Medlock.

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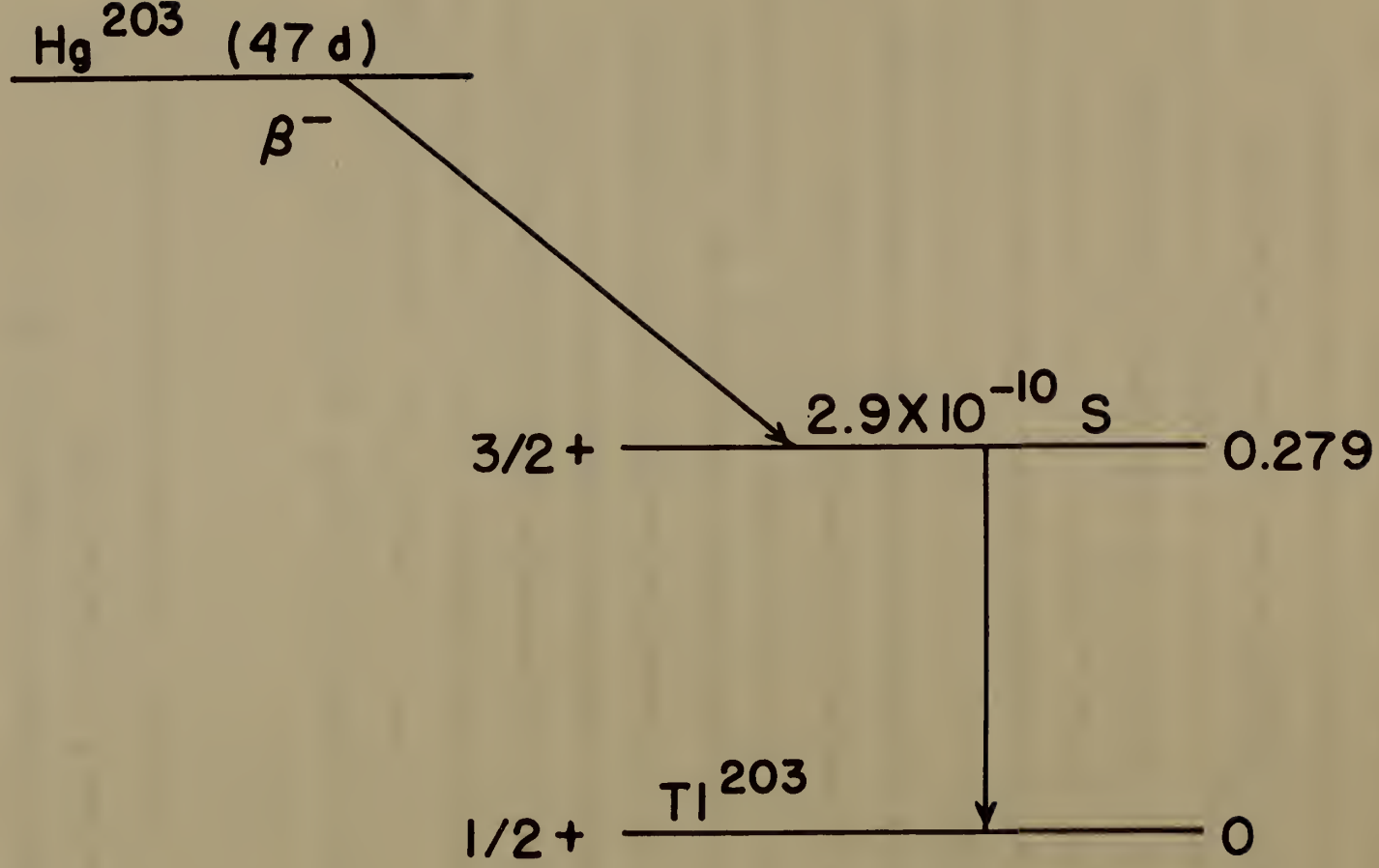



Figure I. Mercury-203 decay scheme. (10)

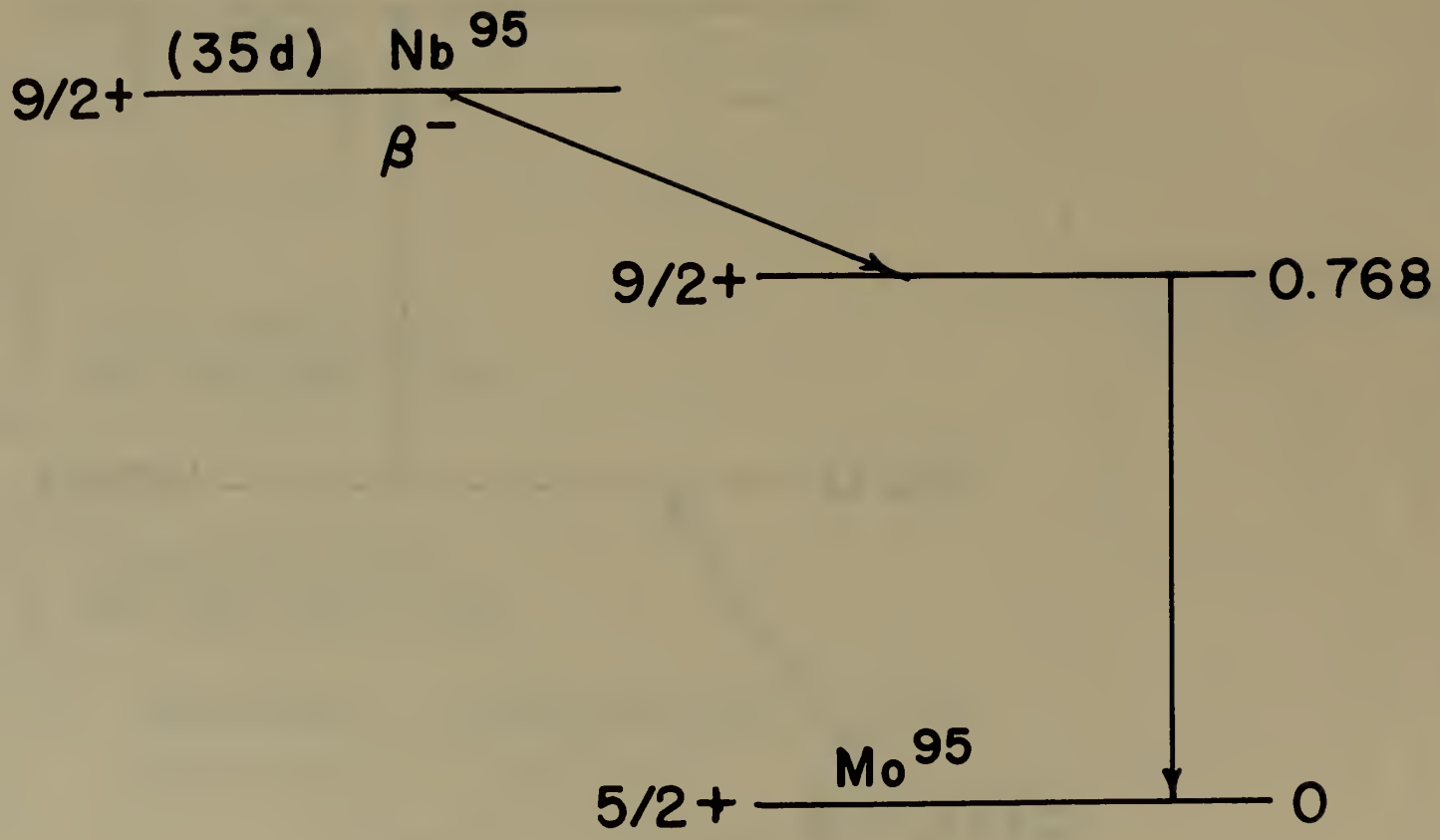



Figure 11. Niobium-95 decay scheme. (11)

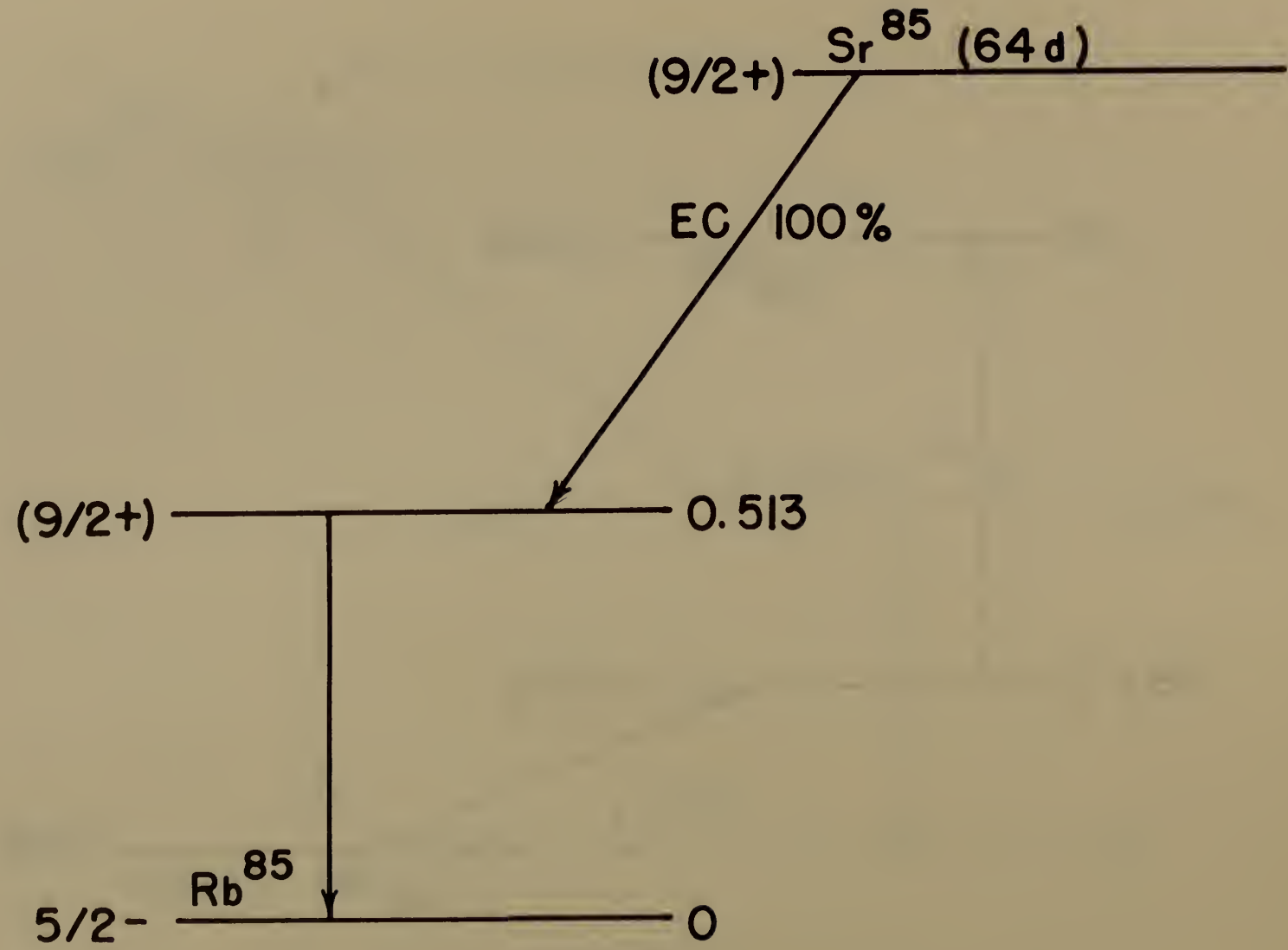



Figure III. Strontium-85 decay scheme. (12)

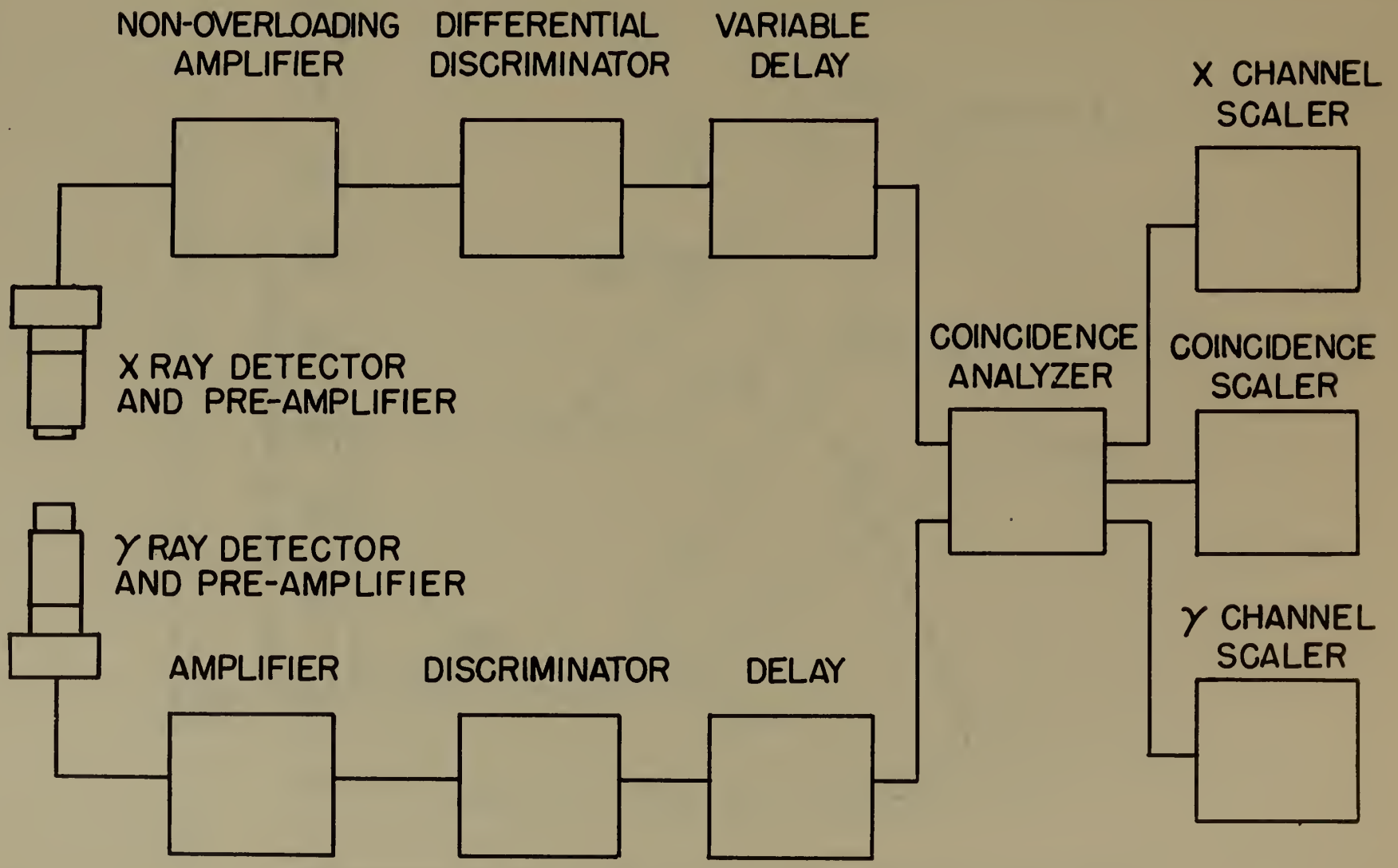



Figure IV. Diagram of the experimental arrangement for the half-life measurement of the isomeric state of Rubidium-85.

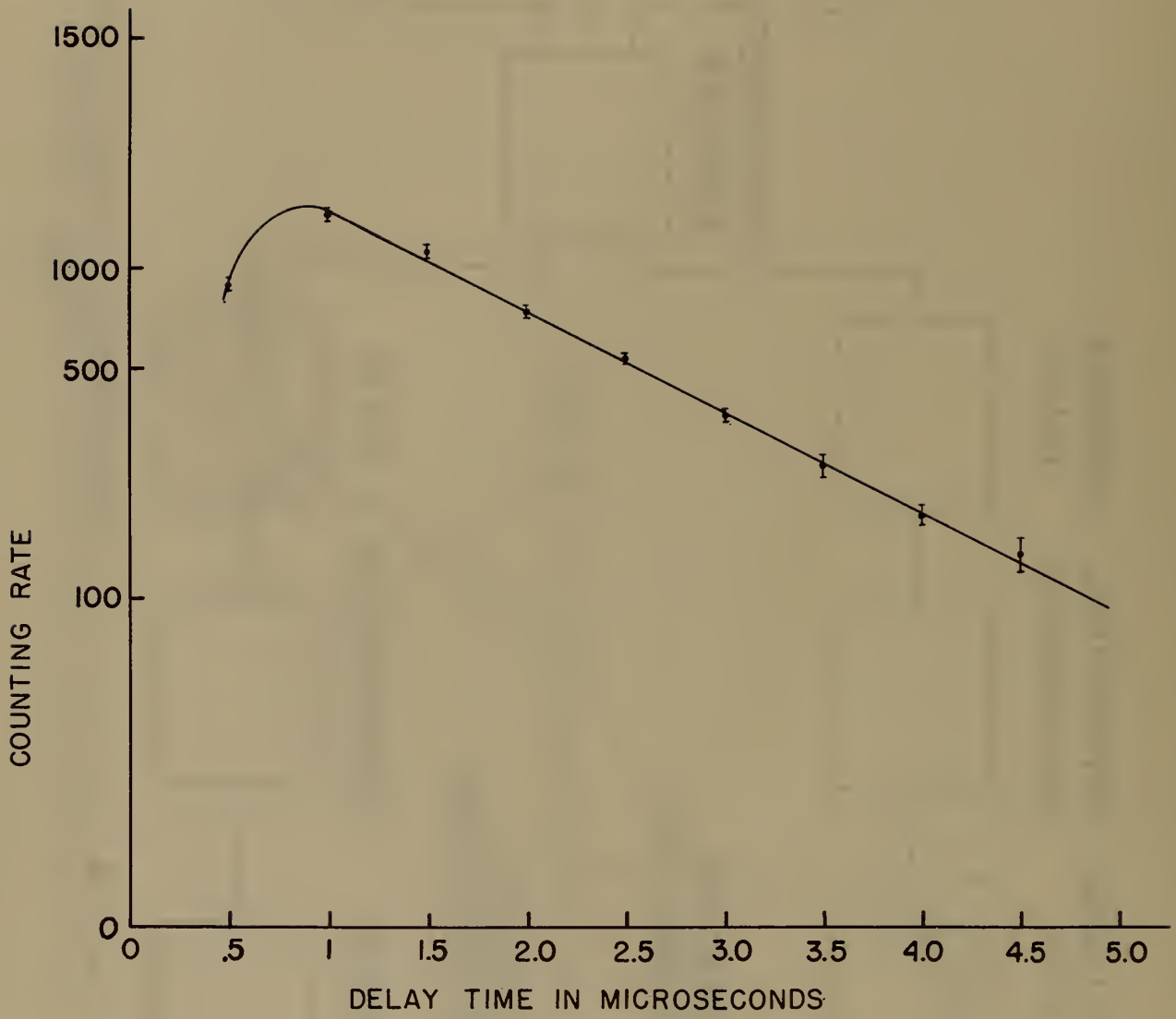



Figure V. Decay curve of the isomeric state of Rubidium-85.

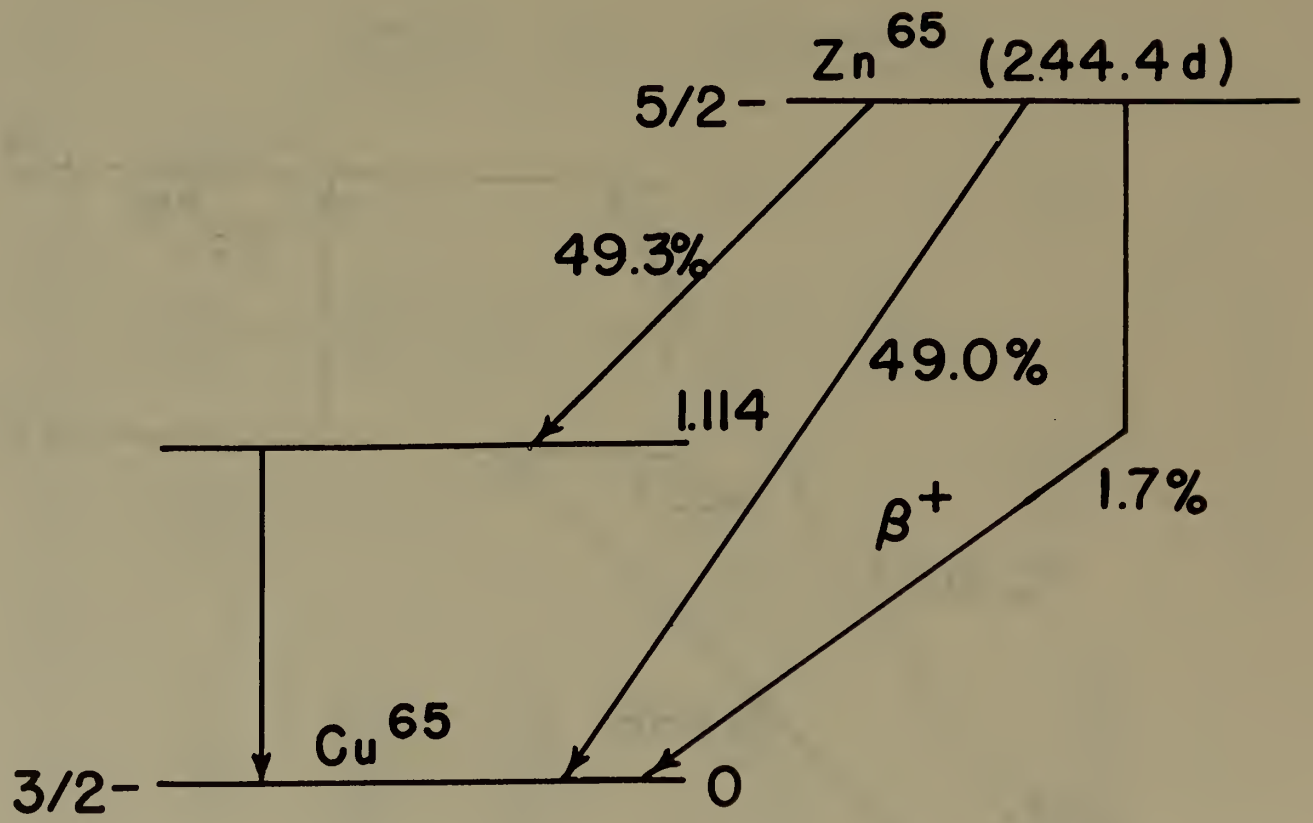



Figure VI. Zinc-65 decay scheme. (13)

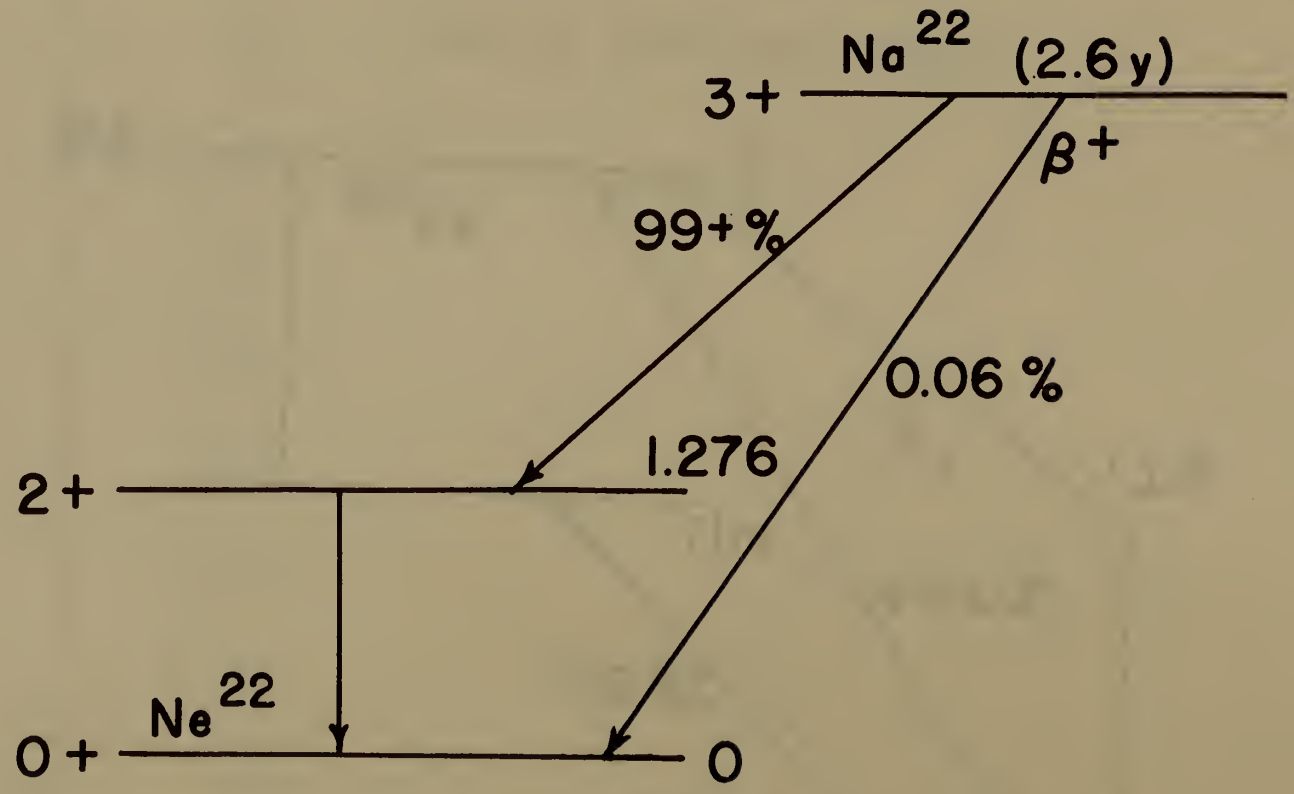



Figure VII. Sodium-22 decay scheme. (14)

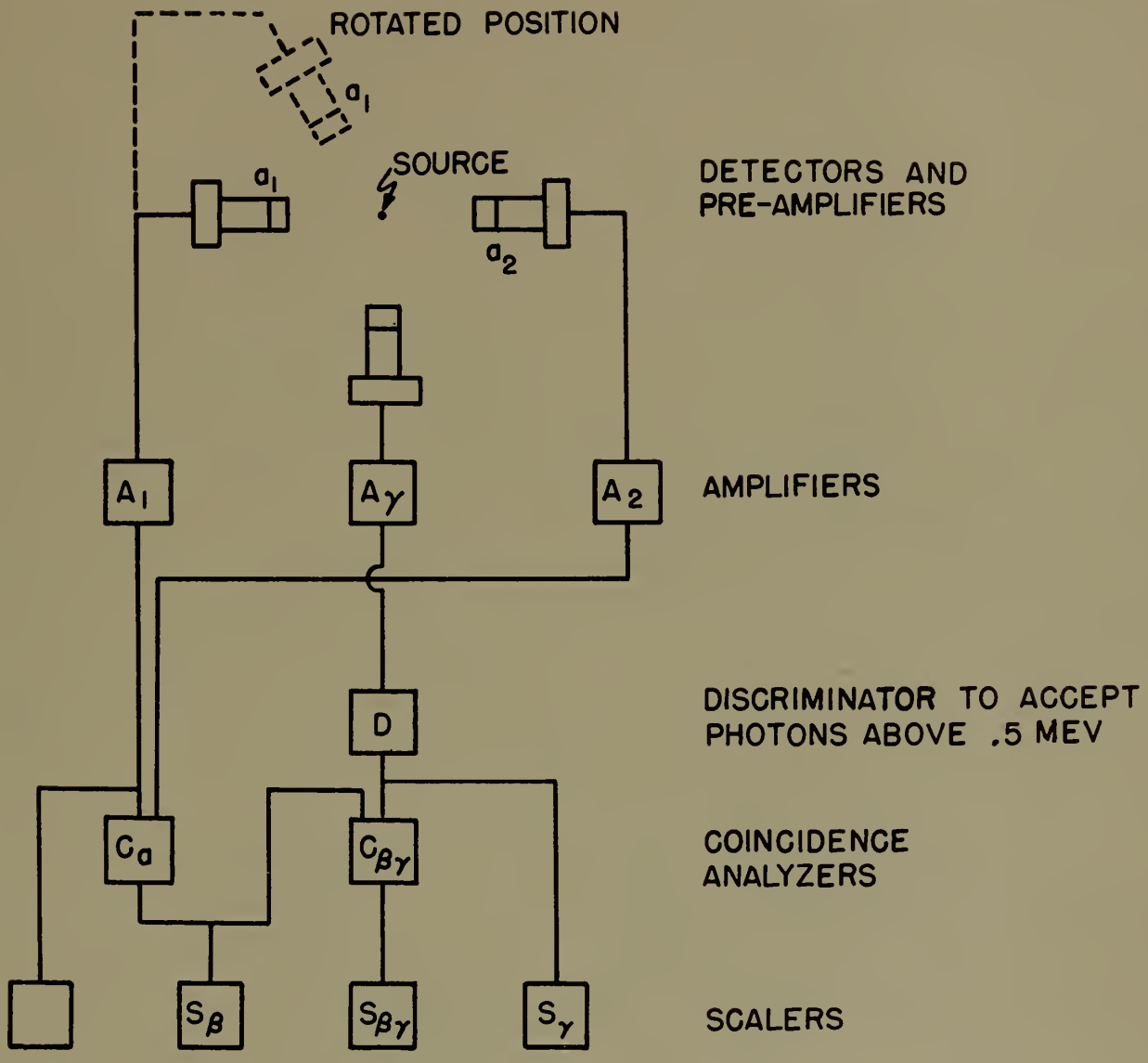
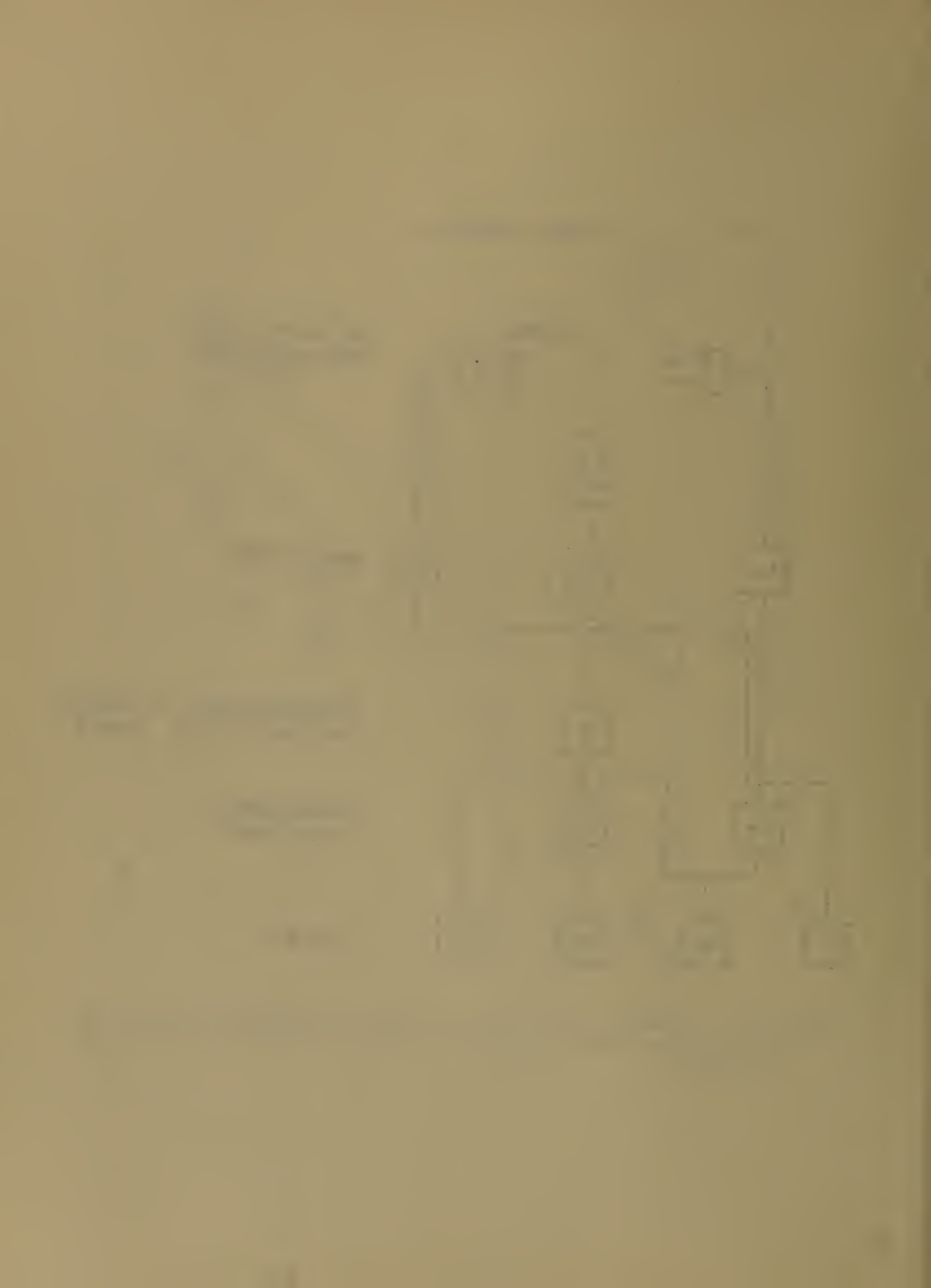



Figure VIII. Schematic diagram of the triple coincidence apparatus for the determination of the absolute disintegration rate of Sodium-22 sources.



U.S. DEPARTMENT OF COMMERCE
Frederick H. Mueller, *Secretary*

NATIONAL BUREAU OF STANDARDS
A. V. Astin, *Director*

# THE NATIONAL BUREAU OF STANDARDS

The scope of activities of the National Bureau of Standards at its major laboratories in Washington, D.C., and Boulder, Colorado, is suggested in the following listing of the divisions and sections engaged in technical work. In general, each section carries out specialized research, development, and engineering in the field indicated by its title. A brief description of the activities, and of the resultant publications, appears on the inside of the front cover.

## WASHINGTON, D.C.

**Electricity and Electronics.** Resistance and Reactance. Electron Devices. Electrical Instruments. Magnetic Measurements. Dielectrics. Engineering Electronics. Electronic Instrumentation. Electrochemistry.

**Optics and Metrology.** Photometry and Colorimetry. Photographic Technology. Length. Engineering Metrology.

**Heat.** Temperature Physics. Thermodynamics. Cryogenic Physics. Rheology. Molecular Kinetics. Free Radicals Research.

**Atomic and Radiation Physics.** Spectroscopy. Radiometry. Mass Spectrometry. Solid State Physics. Electron Physics. Atomic Physics. Neutron Physics. Radiation Theory. Radioactivity. X-rays. High Energy Radiation. Nucleonic Instrumentation. Radiological Equipment.

**Chemistry.** Organic Coatings. Surface Chemistry. Organic Chemistry. Analytical Chemistry. Inorganic Chemistry. Electrodeposition. Molecular Structure and Properties of Gases. Physical Chemistry. Thermochemistry. Spectrochemistry. Pure Substances.

**Mechanics.** Sound. Mechanical Instruments. Fluid Mechanics. Engineering Mechanics. Mass and Scale. Capacity, Density, and Fluid Meters. Combustion Controls.

**Organic and Fibrous Materials.** Rubber. Textiles. Paper. Leather. Testing and Specifications. Polymer Structure. Plastics. Dental Research.

**Metallurgy.** Thermal Metallurgy. Chemical Metallurgy. Mechanical Metallurgy. Corrosion. Metal Physics.

**Mineral Products.** Engineering Ceramics. Glass. Refractories. Enameled Metals. Constitution and Microstructure.

**Building Technology.** Structural Engineering. Fire Protection. Air Conditioning, Heating, and Refrigeration. Floor, Roof, and Wall Coverings. Codes and Safety Standards. Heat Transfer. Concreting Materials.

**Applied Mathematics.** Numerical Analysis. Computation. Statistical Engineering. Mathematical Physics.

**Data Processing Systems.** SEAC Engineering Group. Components and Techniques. Digital Circuitry. Digital Systems. Analog Systems. Application Engineering.

• Office of Basic Instrumentation.

• Office of Weights and Measures.

## BOULDER, COLORADO

**Cryogenic Engineering.** Cryogenic Equipment. Cryogenic Processes. Properties of Materials. Gas Liquefaction.

**Radio Propagation Physics.** Upper Atmosphere Research. Ionospheric Research. Regular Propagation Services. Sun-Earth Relationships. VHF Research. Radio Warning Services. Airglow and Aurora. Radio Astronomy and Arctic Propagation.

**Radio Propagation Engineering.** Data Reduction Instrumentation. Modulation Research. Radio Noise. Tropospheric Measurements. Tropospheric Analysis. Propagation Obstacles Engineering. Radio-Meteorology. Lower Atmosphere Physics.

**Radio Standards.** High Frequency Electrical Standards. Radio Broadcast Service. High Frequency Impedance Standards. Electronic Calibration Center. Microwave Physics. Microwave Circuit Standards.

**Radio Communication and Systems.** Low Frequency and Very Low Frequency Research. High Frequency and Very High Frequency Research. Ultra High Frequency and Super High Frequency Research. Modulation Research. Antenna Research. Navigation Systems. Systems Analysis. Field Operations.

NBS